
Mean-Field limit of the Bose-Hubbard model in high dimension

Denis Périce & Shahnaz Farhat, joint work with Sören Petrat

Les Houches Summer School 2026: Quantum Theory on All Scales - August 19th 2026



August 12th 2026 solar eclipse from the Prarion summit.

- [1] S.Farhat D.Périce S.Petrat. "Mean-Field Dynamics of the Bose-Hubbard Model in High Dimension". In: *Annales Henri Poincaré* (to appear).
- [2] S.Farhat D.Périce S.Petrat. "Ground state energy of the Bose-Hubbard model with large coordination number with a polaron-type quantum de Finetti theorem". In: *arXiv:2603.29562* (2026).

Convergence of the ground state energy

d-dimensional cubic lattice: of length $L \in \mathbb{N}^*$: $\Lambda_d := (\mathbb{Z}/L\mathbb{Z})^d$.

Bose-Hubbard Hamiltonian acting on $\mathcal{F}_d := \ell^2(\mathbb{C})^{\otimes |\Lambda_d|}$:

$$H_d := -\frac{J}{2d} \overbrace{\sum_{\substack{x,y \in \Lambda_d \\ |x-y|=1}}^{\mathcal{O}(2d|\Lambda_d|)} a_x^\dagger a_y + (J - \mu) \sum_{x \in \Lambda_d} \mathcal{N}_x + \frac{U}{2} \sum_{x \in \Lambda_d} \mathcal{N}_x (\mathcal{N}_x - 1).$$

Mean-field Hamiltonian given $\varphi \in \ell^2(\mathbb{C})$:

$$h_\varphi := -J(\overline{\alpha_\varphi} a + \alpha_\varphi a^\dagger - |\alpha_\varphi|^2) + (J - \mu)\mathcal{N} + \frac{U}{2}\mathcal{N}(\mathcal{N} - 1),$$

of order parameter $\alpha_\varphi := \langle \varphi, a\varphi \rangle$.

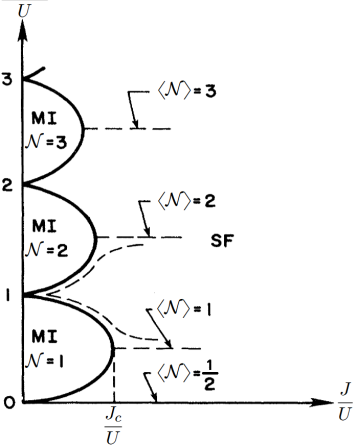
Theorem: *S.Farhat & D.P. & S.Petrat 2026 [2]*

If $U > 0$ and $J \geqslant 0$, then

$$\inf_{\psi \in \mathcal{F}_d, \|\psi\|_{\mathcal{F}_d}=1} \frac{\langle \psi, H_d \psi \rangle}{|\Lambda_d|} \xrightarrow{d \rightarrow \infty} \inf_{\varphi \in \ell^2(\mathbb{C}), \|\varphi\|_{\ell^2}=1} \langle \varphi, h_\varphi \varphi \rangle.$$

Phase transition: for a mean-field energy minimizer φ_0 ,

- Mott Insulator: $\langle \varphi_0, a\varphi_0 \rangle = 0$.
- Superfluid: $\langle \varphi_0, a\varphi_0 \rangle > 0$.



Convergence of the dynamics

Schrödinger equation: Let $\Psi_d \in \mathcal{F}_d$ solve

$$i\partial_t \Psi_d(t) = H_d \Psi_d(t), \quad H_d := -\frac{J}{2d} \sum_{\substack{x,y \in \Lambda_d \\ |x-y|=1}} a_x^\dagger a_y + (J - \mu) \sum_{x \in \Lambda_d} \mathcal{N}_x + \frac{U}{2} \sum_{x \in \Lambda_d} \mathcal{N}_x(\mathcal{N}_x - 1). \quad (\text{BH})$$

Mean-field dynamics: Let $\varphi \in \ell^2(\mathbb{C})$ solve

$$i\partial_t \varphi(t) = h_\varphi(t) \varphi(t), \quad h_\varphi := -J(\overline{\alpha_\varphi} a + \alpha_\varphi a^\dagger - |\alpha_\varphi|^2) + (J - \mu)\mathcal{N} + \frac{U}{2}\mathcal{N}(\mathcal{N} - 1), \quad \alpha_\varphi = \langle \varphi, a\varphi \rangle. \quad (\text{MF})$$

Reduced density matrix: Let $\gamma_d := |\Psi_d\rangle\langle\Psi_d|$ and define the one-lattice site reduced density matrix

$$\gamma_d^{(1)} := \frac{1}{|\Lambda|} \sum_{x \in \Lambda} \text{Tr}_{\Lambda \setminus \{x\}} (\gamma_d), \quad p_\varphi := |\varphi\rangle\langle\varphi|, \quad q_\varphi := 1 - p_\varphi.$$

Theorem: *S.F. & D. Périce & S. Petrat 2025 [1]*

We assume:

- Ψ_d solves (BH) with $\|\Psi_d(0)\|_{\mathcal{F}_d} = 1$,
- φ solves (MF) with $\|\varphi(0)\|_{\ell^2} = 1$,
- $U > 0$ and there is $C > 0$ such that

$$\text{Tr}(p_\varphi(0)\mathcal{N}^4) \leq C.$$

Then, there exists $C(J, \mu, U) > 0$ such that for all $t \in \mathbb{R}_+$,

$$\|\gamma_d^{(1)}(t) - p_\varphi(t)\|_{\mathcal{L}^1} \leq C(J, \mu, U) e^{C(J, \mu, U)(1+t^7)} \left(\text{Tr} \left[\gamma_d^{(1)}(0) \left(q_\varphi(0) + q_\varphi(0) \mathcal{N}^2 q_\varphi(0) \right) \right] + \frac{1}{d} \right)^{1/2}.$$